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C. Lathuillère, M. Menvielle, J. Lilensten, T. Amari, S. M. Radicella. From the Sun's atmosphere to the Earth's atmosphere: an overview of scientific models available for space weather developments. *Annales Geophysicae*, 2002, 20 (7), pp.1081-1104. 10.5194/angeo-20-1081-2002 . hal-00329258

**HAL Id: hal-00329258**

**<https://hal.science/hal-00329258>**

Submitted on 1 Jan 2002

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# From the Sun's atmosphere to the Earth's atmosphere: an overview of scientific models available for space weather developments

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Received: 21 August 2001 – Revised: 27 March 2002 – Accepted: 3 April 2002

**Abstract.** Space weather aims at setting operational numerical tools in order to nowcast, forecast and quantify the solar activity events, the magnetosphere, ionosphere and thermosphere responses and the consequences on our technological societies. These tools can be divided in two parts. The first has a geophysical base (Sun, interplanetary medium, magnetosphere, atmosphere). The second concerns technological applications (telecommunications, spacecraft orbits, power plants ...). In this paper, we aim at giving an overview of the models that belong to the first class (geophysics) that might serve in the future as a basis for building global operational codes. For each model, we consider the physics underneath, the input and output parameters, and whether it is already operational, whether it may become operational in the near future, or if it is an academic research tool. Relevant references are given in order to serve as a starting point for further readings.

**Key words.** Interplanetary physics (general or miscellaneous), Ionosphere (modelling and forecasting), Magnetospheric physics (general or miscellaneous)

## 1 Introduction

Reviewing the current state of scientific models available for space weather developments is actually an extremely ambitious task. It covers a broad range of topics – from the Sun to the Earth – and of techniques and methods – from the ideal Magnetohydrodynamics (MHD) to particle behaviour modelling. To do a thorough job would have required a book. However, the range of topics involved in space weather makes it necessary to have introductory documents that pro-

vide the non-specialist with basic features and relevant references on the various issues.

The present work aims at providing an overview of the models which might serve as a basis to build global operational codes for practical space weather prediction. Our objective is to provide the reader with basic features on the various issues, and not to deal with each of them in any sufficient or satisfying detail for a specialist. It is also to give relevant references that can serve as a starting point for those who are interested in specific topics.

This paper has been written from the work package entitled “Space Weather Parameters”, prepared as part of a study conducted by the consortium lead by ALCATEL Space & LPCE (France) in the frame of an ESA contract aiming at the definition of an European Space Weather Programme. All work packages have received the contribution from the different consortium members (see Annex 1 for a presentation of the consortium). A parallel study has been undertaken by another consortium led by the Rutherford Appleton Laboratories (UK). Among the tasks of the consortium was to define the parameters and the associated measurements, which are necessary to describe and monitor the Sun-Earth system in the space weather context.

Space weather can be considered from two different standpoints, the numerical modelling of the whole system on the one hand, and the needs of the users, i.e. monitoring and forecasting the values of some specific quantities, on the other hand. Our approach has been to start with a review of the modelling effort, which is presented in this paper.

The driving source of the Solar Terrestrial phenomena is the Sun's magnetic field. The related magnetic energy is converted into thermal and kinetic energy, and gives rise to electromagnetic and particle solar radiations. The plasma that is ejected from the upper solar atmosphere is accelerated to supersonic speeds, and the Sun's magnetic field lines are frozen

in the plasma. The Earth's magnetic field behaves as a solid obstacle for the plasma. It gives rise to the magnetospheric cavity around which the solar wind flows, with energy and material exchange through the magnetosphere boundary. Finally, both magnetospheric plasma and solar electromagnetic radiations interact with the Earth's atmosphere, giving rise to the ionosphere and thermosphere. This very short summary shows that the entire Solar-Terrestrial system can be described and understood in terms of a succession of sub-systems that exchange material and energy: the Sun's atmosphere, the interplanetary medium, the magnetosphere, and finally, the ionosphere-thermosphere system. Figure 1 shows a schematic description of the solar-terrestrial phenomena in terms of the interaction between these sub-systems.

The elucidation of the physical processes involved in each sub-system, and in the interaction between any two of them, is an important key in any space weather program. Our understanding of the whole Sun-Earth system is, however, still incomplete, although research and operational models are available. There exist research models for each sub-system, but their maturity is very different, depending on the complexity of physical phenomena and on the available observations. Operational models have already been developed only for some sub-systems. Our aim is not to present an exhaustive list of models but rather to give a comprehensive picture of the modelling effort of the different scientific communities concerned by space weather.

We have chosen to organise our presentation by classifying for each sub-system the models in two main classes:

- the empirical models that characterise the relations between relevant parameters from available observations;
- the physical models that describe a given sub-system on the basis of quantitative laws between relevant parameters. Contrary to the empirical models, the relations between parameters are expressed in terms of a priori known physical laws, with relevant simplifications, if any. They are expressed as a set of partial derivative equations, with coefficients to be estimated.

Recently developed space weather models based upon artificial techniques, such as neural networks genetic algorithms and expert systems, are not presented here, nor are the technological models which estimate the specified effects of our environment on a given “system”, for example, those dealing with radiation doses, spacecraft charging, proton fluences or atmospheric drag.

In this paper, we consider successively the different sub-systems that are involved in solar-terrestrial relations: the Sun, the interplanetary medium, the magnetosphere, and the ionosphere-thermosphere system. The last section of the paper presents a model synthesis, by means of tables that outline their main characteristics, including their input and output parameters.

## 2 The Sun

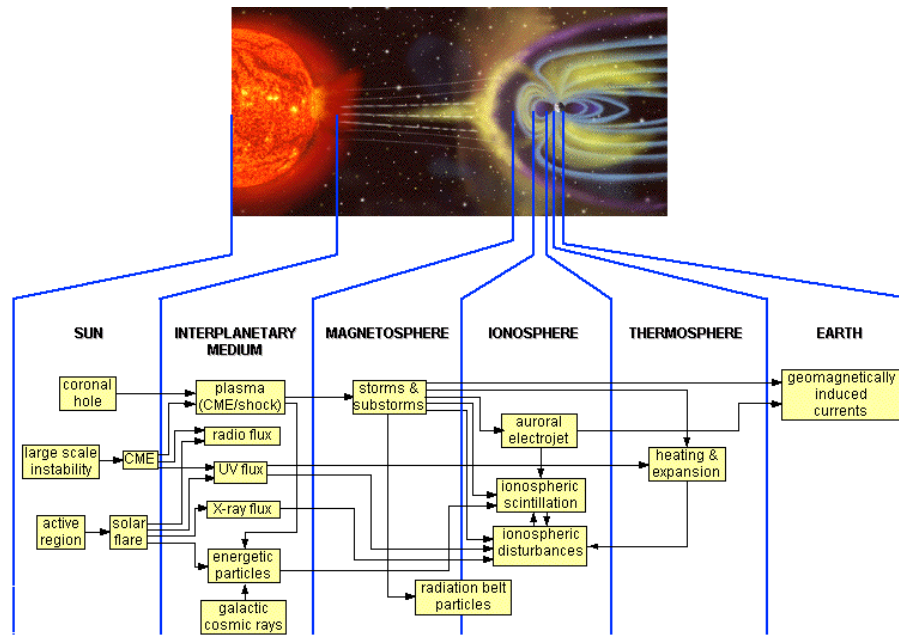
Most of the structures and phenomena present in the solar atmosphere – in particular eruptive phenomena, such as flares and coronal mass ejections (CMEs) – result from the presence of a dominant magnetic field. Eruptive events correspond to a liberation of magnetic energy stored in the solar corona. This energy is then converted into:

- heating of the environment associated with UV/EUV and X beams;
- particle accelerations (electrons and ions) associated with an X emission when those particles interact with the environment;
- movement of matter.

Solar flares correspond to localised phenomena covering at most a few percent of the surface of the Sun, while CMEs are larger scale phenomena that can involve a non-negligible part of the Sun's global configuration. Those may have different origins and associated events, such as global magnetic non-equilibrium or prominence disruption.

The global solar atmospheric models are developed mainly in the framework of Magnetohydrodynamics (MHD). The relevant set of equations describes the interaction of ionised coronal plasma with the coronal magnetic field in the presence of the plasma pressure and gravity forces. These models may be presented in the context of space weather, although they are still essentially used for theoretical purposes. The associated numerical codes are mostly research tools due to the actual state of the art in solar MHD modelling. We present thereafter two main classes of complementary models used for different purposes: the static models for equilibrium reconstruction of the solar coronal magnetic field and the dynamic models to describe its evolution:

- the first class of models arises from the impossibility of measuring the coronal magnetic field. The structure of the active regions should be estimated before an eruptive event, in order to determine the intrinsic properties of the magnetic configuration. Thus, one has to reconstruct the coronal magnetic field and, therefore, to solve the equations of the solar atmospheric physics when the boundary conditions are the values of the magnetic field measured in the colder photosphere by vector magnetographs;
- the second class of models aims at studying the dynamical evolution of the active regions. The energy storage and energy release in these regions, as well as their stability, are described from the evolution of the magnetic configurations which are constrained by a driver whose origin may be sub-photospheric (emerging flux), photospheric (boundary motions) or coronal (interaction with other active regions). These models solve – to a certain extent – the full MHD equations.



**Fig. 1.** Main physical processes that act on space weather (Dudok de Wit, personal communication, 2001).

Both classes of models are complemented by a whole set of boundary conditions which then define a set of boundary value problems.

## 2.1 Reconstruction and study of the active region static structures

Equilibrium reconstruction of the coronal magnetic field above active regions using photospheric magnetic data has been the subject of numerous studies since the first attempts (Schmidt, 1964). They may range from observational problems, such as those related to the  $180^\circ$  ambiguity resolution, which remains on the transverse component of the photospheric magnetic field, to crucial theoretical problems related to the nature and the determination of the correct type of boundary conditions that have to be used in order to avoid an ill-posed problem, as was the case for a long time. Details and numerous relevant aspects of these studies are discussed in Amari and Démoulin (1992) and Amari et al. (1997).

In an active region, prior to any eruptive event, typical time variations are so small that one can then solve the MHD equations under static hypothesis. The coronal magnetic pressure is much larger than the gas one, and gravity can be neglected outside prominences. This is the so-called force-free approximation. This is equivalent assuming that at any point the current density and the magnetic field are parallel. The proportionality factor  $\alpha(\mathbf{r})$  and the magnetic field are determined by the boundary conditions defining the boundary value problem.

In the regions of higher density, the plasma pressure gradients and the gravity have been included. This can be done self-consistently in full MHD methods or more qualitatively by seeking magnetohydrostatic solutions of the MHD equations.

### 2.1.1 Current free (potential) model

The simplest physical approximation is the so-called “current free” approximation ( $\alpha = 0$ ), which only requires the longitudinal photospheric component of the magnetic field as a boundary condition. It was first considered by Schmidt (1964) and is currently routinely used in most of the terrestrial solar physics centres, on the basis of observations (Sakurai, 1989). It is also used as the initial conditions for fundamental MHD studies dealing with synthetic problems. This kind of reconstruction is performed following a Green function approach or using a Laplace solver to compute a potential scalar function and the associated magnetic field (Priest, 1982).

### 2.1.2 Linear force-free model

The zero-current approximation does not apply to many active regions that have a magnetic energy above the minimum energy, which corresponds to the current free field for the same distribution of the vertical distribution of the photospheric normal magnetic field. The first step towards a more realistic modelling consists of considering a nonzero, but constant  $\alpha$  that allows one to introduce coronal electric currents. Different methods exist that use the longitudinal component of the magnetic field. They are based on Green functions (Chiu and Hilton, 1977) or Fourier transform (Alissandrakis, 1981), with the latter currently used (Démoulin et al., 1997). Following the general approach of Low (1992), this linear constant- $\alpha$  Fourier method has been extended to take into account gravity and pressure forces by means of a linear computational program (Démoulin et al., 1997; Aulanier et al., 1998).

Although they suffer from several limitations (see below), linear models can be used for some classes of problems (topology, prominence) in moderately sheared structures. Their main advantage remains the weak computational resources they need and their computational speed for the current free approximation, which explains why they could be used routinely in a space-weather program.

### 2.1.3 Nonlinear force-free model

As for the reconstruction without current, several limitations exist regarding the use of the linear reconstruction model:

- the solutions correspond to a minimisation problem for the energy, constrained by the normal component of the magnetic field and by the total magnetic helicity, inducing a limited amount of available magnetic free energy;
- the electric currents cannot be locally intense, while the observations show very clearly important localised shear concentrated along the inversion line of the normal component of the photospheric magnetic field.

The only hope of incorporating large localised electric currents is to assume that the configuration is in a nonlinear force-free state, that is to say that  $\alpha$  (which now depends on the position  $\mathbf{r}$ ) is also an unknown of the problem. Different types of boundary conditions define different boundary value problems and codes:

- a first relatively natural method consists of imposing the three components of the magnetic field  $\mathbf{B}$  measured at the photospheric level. The problem amounts to progressively extrapolating the data step by step toward the corona. This is the vertical integration method introduced by Wu et al. (1990). However, this method is based on a mathematical formulation associated with an ill-posed boundary value problem. This lead to an exponential divergence, which limits the reconstruction at low altitudes. Some methods have been proposed in order to avoid this divergence (Cuperman et al., 1990; Demoulin et al., 1992). Other attempts have been made to regularise the linear version of this method (Amari et al., 1998);
- a second class of reconstruction methods is based on a well-posed formulation, which corresponds to observed boundary conditions that imply:
  - the normal component of the photospheric magnetic field (in the local system of reference associated with the Sun);
  - $\alpha$  at the photospheric level, when the sign of the normal component  $B_z$  of the measured magnetic field has a given a priori value: as an example,  $\alpha$  is set in the areas where  $B_z$  is positive and then computed through transport along the magnetic field lines to the photospheric zones where  $B_z$  is negative.

This kind of mathematical formulation was introduced by Grad and Rubin (1958) and consists of a decomposition of the nonlinear problem in two sets of elliptical (for the magnetic field) and hyperbolic (for  $\alpha$ ) problems. It gives rise to two types of numerical codes, which correspond to the Lagrangian and Eulerian approaches:

- Sakurai (1981) uses a Lagrangian approach for each field line. He injects progressively the photospheric electric current on each field line;
- Amari et al. (1997) use a global approach in order to solve the elliptical and hyperbolic problems. The corresponding codes with their different versions constitute the EXTRAPOL code.

These two codes, localised in Japan and in France, are used to reconstruct the configurations of the active regions coming from MITAKA data for the Japanese version and from Hawaii (HSP, IVM) and Boulder (ASP) for the French one. The local Lagrangian approach (Japan) seems to imply a more important limit for the reproducible maximum shearing than the Eulerian approach (France);

- a third class of methods is based on solving the MHD equations, injecting the vertical component of the electric current and imposing the normal component of the magnetic fields as photospheric boundary conditions. The system relaxes toward a state which corresponds to the observed values. This method is called the “resistive-relaxation” method. It was introduced by Mikic and Linker (1994) and was also implemented in the French METEOSOL MHD-research code. Although the mathematical justification for this method is not yet clear, the corresponding codes have been used on real cases with relatively good success;
- another method is still at a very experimental state of development. It is the “weighted residue method” (Pridmore-Brown, 1981). This method minimises two residuals: one is associated with the magnetic force, and the second one is associated with the difference between the directions of the photospheric transverse components of the computed and observed magnetic fields. This method has been used only in theoretical situations corresponding to fictive periodic data;
- finally, a method called “constraint and relaxation method” (Roumeliotis, 1996) consists of two repeated steps:
  - impose on the vector potential the minimisation of the difference between the computed and observed transverse components;
  - relax toward an equilibrium state, using a relaxation code which incompletely solves the MHD equations,
  - then, back to the first step.

**Table 1.** A synthesis of MHD codes for the evolution of the solar magnetic configurations

Laboratories	Characteristics	Aims
San Diego (USA) (Mikic and Linker, 1994)	2.5D, 3D Finite differences Semi-implicit	coronal MHD
NRL et NASA (Wash. DC) (De Vore et al., 2000)	2.5D, 3D Spectral + finite volumes Explicit Boundary conditions ( . . . )	Coronal and internal MHD
Japan NAO/Mikata (Matsumoto et al., 1993)	2.5D, 3D Explicit artificial viscosity anomalous	Chromospheric and internal MHD
Observatoire de Meudon Ecole Polytechnique (France) (Amari et al., 2000)	2.5D, 3D differences and finite volumes (semi) and implicit	coronal MHD
Strasbourg (France) (Baty and Heyvaerts, 1996)	cylindrical Boundary conditions	coronal MHD stability
Pisa / Firenze (Italy) (Lionello et al., 1998)	cylindrical (3D) San Diego 2D reduced MHD	coronal MHD turbulence
Nice (France) (Galtier et al., 1997)	1D, 2D Finite differences, Spectral	Turbulence Intermittence
Argentina (Dmitruk et al., 1998)	Reduced MHD (2D+) Cartesian Spectral (Fourrier)	Turbulence Flare-heating
University of Michigan (USA) (Israelevich et al., 2001)	2D, 2.5D, 3D (?) Roe scheme Finite volumes	Comets – wind astrophysics
NCSA (Illinois) (Stone and Norman, 1992)	2.5D, (3D ?) Finite volumes Van Leer, PPM	Astrophysics laboratories
DAEC + DESPA (France) (Grappin et al., 2000)	cylindrical /(Spherical 2.5D) Z-periodic Spectral – Finite differences	Dynamo + wind laboratories
Chicago (USA) (Fryxell et al., 2000)	cartesian	Astrophysics laboratories

This method has been tested with encouraging results on a theoretical case, as well as on an observed active region.

## 2.2 Evolution of the magnetic configurations

As far as the evolution of the magnetic configurations is concerned, the MHD codes belong to the domain of research even more than the reconstruction codes. It concerns the fundamental numerical research for the codes themselves: boundary conditions, geometry, temporal scheme, mesh definition, etc., but also the fundamental solar physics. The objective is to elucidate the fundamental mechanisms that govern the eruptive phenomena: flux emergence, energy storage,

influence of the photospheric movements such as differential rotations or magnetic reconnection, etc.

The existing codes depend on the different classes of problems they address: there is no universal MHD code, in the same way as there is no universal telescope. However, it will still be a long time before the MHD codes can be used for space weather purposes, in order to test the stability of the magnetic configurations in active regions. These regions may be first reconstructed by using an equilibrium reconstruction code. This is the approach followed by some groups.

The existing world wide codes can be distinguished by the following characteristics:

- *dimension*: 2D, 2.5D, where the vectors have 3 components but only depend on 2 space variables, 3D;
- *geometry*: spherical for the global models, Cartesian for the description of the active regions at local scales, cylindrical for more specific problems, such as the stability of the coronal loops or of the coronal heating;
- *numerical scheme* and *space digitisation*: finite elements, finite volumes, spectral methods, collocation, etc.
- *physical approach*: heating, resistivity, solar wind inclusion facing the particularly important and difficult problem of breakdown of the MHD approximation.

Table 1 only attempts to synthesize these different approaches worldwide. We are aware that on-going work may give rise to new codes in the near future. We only show here the codes and the associated developing groups and give only one reference per group, as a starting point. The more numerous groups that only use the codes are excluded from the table.

### 3 The interplanetary medium

The solar wind is a prominent part of the Sun-Earth system that is difficult to model because observations are difficult and often indirect. The solar wind corresponds to the outflow from the upward solar atmosphere of the coronal plasma, heated up to several million K. Its speed ranges between 250 and 950 km.s<sup>-1</sup> at the Earth's orbit (1 AU). It is structured by the Sun's magnetic field lines frozen in the flow. The solar wind interacts in the interplanetary medium with galactic cosmic rays. Since the solar wind expansion does not occur in a vacuum but in an interstellar space filled with neutral and ionised gases, galactic cosmic rays and galactic magnetic fields, the region where the solar wind dominates is limited by the heliospheric boundary – a region where the energy densities of the solar wind and the interstellar space are equal. This boundary is located between about 50 and 100 AU.

#### 3.1 Solar wind modelling

The main structures or regions present in the solar wind are:

- the high velocity solar wind, associated with the coronal holes;
- the low velocity solar wind;
- the regions of interactions between high and low velocity solar winds;
- the heliospheric current sheet;
- the signatures in the interplanetary medium of the CMEs, associated or not associated with shocks. In practice, these signatures are observed at large distances from the Sun near the Earth or the Lagrange point L1.

The influence of the solar wind results from both its very structure and its interaction with the propagation of solar perturbations, such as energetic particles resulting from solar eruptions or interplanetary shocks or CMEs. The structure of the interplanetary magnetic field lines, which drives the particle propagation, may, in given situations, dramatically affects this propagation and, therefore, the resulting effects at Earth. This ambivalent action can be illustrated by examples of major events at the Sun with minor effect's at the Earth, or on the contrary, minor events at the Sun resulting in a major perturbation at the Earth.

Solar wind models have been developed to address many aspects of its structure. Fairly simple models can be developed by treating the solar wind from a hydrodynamic point of view, and thus solving the one-dimensional equations of motion along a single field line (e.g. Parker, 1963). The main aim of such models is to obtain accurate expressions for the plasma density, velocity, and especially temperature at 1 AU. They are also amenable to adding multiple ion species and to including quite detailed calculations of different ionisation states and even the effect of plasma instabilities on the temperatures.

However, from the point of view of space weather, it is the large-scale multi-dimensional MHD models that are the most relevant. One class of these models examines the global structure of the solar wind plasma and magnetic field in the heliosphere. These models are quasi-steady state, but incorporate solar rotation, and a three-dimensional magnetic field. For solar minimum conditions – for which they are best suited – they clearly demonstrate the formation of the familiar co-rotating interaction regions at distances of 1 AU and beyond (Pizzo, 1982, 1991). The Ulysses mission has permitted a comparison of these results with the three-dimensional solar wind structure. The agreement is good, especially in terms of plasma flows upward and downward from the ecliptic plane (Riley et al., 1996).

A second family of models deals with the motion of coronal mass ejections in the solar wind. The goal is, for given plasma and magnetic field conditions at the Sun, to calculate the properties of the CME at 1 AU and beyond. Early models treated the CME as a pressure pulse, but these are now viewed as irrelevant. Other purely hydrodynamic models have been developed by Riley and collaborators during an investigation of so-called over-expanding CMEs. They showed that with a large plasma over-pressure at the Sun, the conditions observed by the Ulysses spacecraft at large distances could be roughly reproduced (Riley et al., 1997).

The most geo-effective type of CMEs are magnetic clouds. Models for these have been developed by a number of groups who have – in general terms – established that magnetic clouds can propagate from the Sun to the Earth while retaining their organised magnetic structure (e.g. Cargill et al., 2000; Odstrcil and Pizzo, 1999; Vandas et al., 1996). The clouds interact with the solar wind by two processes. First, the cloud magnetic field can undergo magnetic reconnection with the solar wind field, leading to its ultimate destruction at large distances. Second, the interaction of the cloud with the

solar wind plasma leads to considerable changes in its shape, as well as leading shock waves. This modelling is plagued by difficulties in initialising the solar wind magnetic field in a way that minimises errors in the  $\text{div}\mathbf{B} = 0$  condition. Proper 3-D modelling actually requires that this problem be resolved. Further discussion of such CME models can be found in Cargill and Schmidt (this issue).

Among the empirical models, one can quote Wang-Sheeley, based upon the discovery of an anti-correlation between the solar wind velocity and the rate of expansion of the magnetic flux tubes (Wang and Sheeley, 1992). The model is based on observations of magnetographs at the Sun's surface and on deduced maps of coronal holes, surface magnetic field sources and solar wind velocity. It forecasts the IMF polarity and the solar wind velocity at the Earth, two parameters that are mandatory for geomagnetic activity forecasting<sup>1</sup>.

### 3.2 Solar wind – cosmic ray interaction

Cosmic rays at the Earth are high energy particles – from 500 MeV to about  $10^{12}$  GeV – insensitive to the magnetosphere state. The main components are protons, electrons and light nuclei. Below about 1 GeV/nucleon, the cosmic ray flux is strongly decreased due to adiabatic deceleration with the solar wind, which results a decreasing flux below this energy at the Earth's orbit. The intensity of the cosmic rays on the Earth maximises when solar activity is minimal (quiet Sun) and minimises vice versa (active Sun), with an average variation in intensity of about 20%.

Models describing the propagation of the cosmic rays through the solar wind are based on a second order equation whose main input parameters are the solar wind characteristics and the interplanetary magnetic field. Fisk et al. (1998) show that most of the observed phenomena can be accounted for within acceptable ranges of these parameters. However, no single model, with a single choice for the input parameters, has been able to account for all the observed features of galactic and anomalous cosmic ray behaviour. The most highly developed models concern the interpretation of long-term variations of the flux of cosmic rays at the Earth.

## 4 The magnetosphere

The magnetosphere is the region of the ionised environment of the Earth where the Earth's magnetic field has a dominant control over the motion of charged particles. The boundary layer between the magnetosphere and the solar wind is the magnetopause, at which the dynamic pressure of the solar wind is balanced by the magnetic pressure of the Earth's magnetic field<sup>2</sup>. The location of the magnetopause obviously

depends on the status of the solar wind. Under typical solar wind conditions, the terrestrial magnetosphere extends up to  $\sim 10$  Earth radii ( $R_E$ ) in the sunward direction and to several hundreds  $R_E$  in the antisunward direction. The solar wind velocity is supersonic at 1 AU, and the magnetosphere behaves as a solid obstacle. Therefore, there is a bow shock upstream from the magnetopause. In the solar direction, the bow shock is located at a distance in the range of 2–4  $R_E$  from the magnetopause. The region between the bow shock and the magnetopause is the magnetosheath, in which the shocked solar wind plasma has a lower velocity and a temperature 5 to 10 times higher than in the solar wind.

The inner magnetosphere is the region of major concern for space weather issues. Three regions have been identified there following different criteria: the plasmasphere, the radiation belts, and the ring current:

1. *the plasmasphere* corresponds to a cold ( $T_e \approx 1$  eV) but dense ( $n_e \approx 5 \times 10^2 \text{ cm}^{-3}$ ) plasma which co-rotates with the Earth. It is a torus-shaped volume in the innermost magnetosphere. The outer boundary of the plasmasphere is the plasmapause, where the density sharply drops down to about  $1 \text{ cm}^{-3}$ . On average, the plasmapause is located about 4  $R_E$  in the equatorial plane, i.e. a McIlwain parameter<sup>3</sup>  $L = 4$ . It may reach out to  $L = 5 - 6$  during periods of magnetospheric quietness and be compressed down to  $L = 3$  during periods of intense magnetospheric activity;
2. *the ring current* refers to “those parts of the particles in the inner magnetosphere which contribute substantially to the total current density” (Hultqvist et al., 1999). It is then defined with regard to its magnetic signature, and corresponds to a toroidal-shaped electric current that flows westwards around the Earth, with variable density at geocentric distances between  $\sim 2 R_E$  and  $\sim 9 R_E$ . During magnetospheric storms, the particles that contribute substantially to the total current density are mainly trapped ions in the medium energy range (few tens of keV to few hundreds of keV) that originate in the solar wind ( $\text{He}^{++}$ ), the plasmasphere and the ionosphere ( $\text{O}^+$ ). During low activity periods, the particles responsible for the ring current are mostly protons. This change in composition impacts the evolution with time of the ring current, because loss mechanisms (wave-particle interactions, charge exchange, Coulomb collisions, etc.) depend on the mass and energy of the particles;
3. *the radiation belts* generally refer by now to the “high energy ions and electrons that can penetrate into spacecraft shielding materials and eventually cause radiation damage to spacecraft instrumentation and to humans”

the Earth's magnetic field at the magnetopause.  $\mu$  is the magnetic permeability of the medium.

<sup>3</sup>For a given geomagnetic field line, the Mac Ilwain parameter  $L$  is the geocentric distance of the point where it crosses the magnetospheric equator. It is expressed in Earth radii ( $R_E$ ).

<sup>1</sup><http://solar.sec.noaa.gov/~narge/>

<sup>2</sup>The dynamic pressure of the solar wind is  $2\rho V^2 \cos^2 \chi$ , where  $\rho$  and  $V$  are the solar wind density and velocity, and  $\chi$  is the angle between the direction of the solar wind velocity and that of the normal to the magnetopause; the magnetic pressure of the Earth's magnetic field is  $B^2/2\mu$ , where  $B$  is the tangential component of



(Hultqvist et al., 1999). The radiation belts are defined with regard to the energy of particles trapped in the geomagnetic field lines, above  $\sim 100$  keV and reaching values of hundreds of MeV. Most of the trapped particles are protons and electrons, giving rise to the proton and electron radiation belts:

- *the proton belt* has generally one maximum, which corresponds to a McIlwain parameter  $L$  that decreases with increasing proton energy. For 1 MeV, the maximum corresponds to  $L \sim 2.5$  (i.e. field lines crossing the equatorial plane at an altitude  $h_{eq}$  of  $\sim 10\,000$  km), while it corresponds to  $L \sim 1.4$  (i.e.  $h_{eq} \sim 2\,500$  km) for 50 MeV. The proton belt is fairly stable. At low altitude, it is modulated by the solar cycle, in particular in the region of the South Atlantic magnetic anomaly: its intensity is maximum during solar minimum, and minimum during solar maximum. Major magnetic storms (e.g. March 1991) may significantly impact the proton belts, in particular for high McIlwain parameters values;
- *the electron belt* has generally two maxima, giving rise to the internal and external electron belts. The internal belt, with a maximum corresponding to 500 keV around  $L \sim 1.4$  (i.e.  $h_{eq} \sim 6\,400$  km), does not significantly vary with time. On the contrary, the external belt, the maximum of which is around  $L \sim 4$  (i.e.  $h_{eq} \sim 19\,000$  km) fluctuates dramatically under the control of strong magnetic storms: short-term variations of the flux, up to 4 orders of magnitude in a few hours, can be observed during periods of intense magnetic activity. The long-term variation of the external belt fluxes is driven by the solar cycle: the annual mean values of the flux are maximum during the descending phase (about 3 years after the maximum), when coronal holes are in good conjunction with the Earth.

It is worth noting that ring current, radiation belts and the plasmasphere partially overlap. For instance, at  $L = 3$ , the density of the cold plasma is about 1000 times higher than the density of energetic protons ( $> 100$  keV), whereas the energy density of energetic protons dominates by a factor of about 1000. It is also worth noting that trapped radiation belts and the ring current are actually closely related, since the major part of the ring current is carried by trapped particles and all the trapped particles contribute to the ring current (Hultqvist et al., 1999). For a thorough description of the present knowledge on these regions of the magnetosphere, refer to the recently published reviews on the ring current (Daglis et al., 1999) and the inner magnetosphere (Hultqvist et al., 1999).

## 4.1 Global magnetosphere modelling

### 4.1.1 Empirical models

Empirical models for the bow shock and magnetopause have been developed for decades, in the case of the Earth, as well as that of other magnetised planets (e.g. Slavin and Holzer, 1981). Shue et al. (1997) recently published a well documented model, based on fresh data. It has a simple functional form driven by two adjustable parameters: the stand-off distance in the solar direction and the tail flaring.

Empirical models for the magnetic field inside the magnetosphere are based on a statistical analysis of the available magnetic field observations, parameterised by geomagnetic indices. They rely basically on a combination of the Earth's planetary magnetic field – usually described by the IGRF model – and external fields estimated from both in situ magnetic field measurements and mathematical modelling of the current systems (e.g. Tsyganenko, 1990; 1995; Hilmer and Voigt, 1995). These models are updated continuously to account for more and more complex processes in the magnetosphere.

At a given time, the dynamics of the magnetosphere depends on both the present solar wind conditions and magnetosphere status. This status, therefore, depends on the past history of the magnetosphere, and models that only use the present magnetic field values as input cannot provide reliable predictions of the magnetospheric state. In order to develop better predictors, indicators of the magnetosphere history should be added to the inputs.

### 4.1.2 MHD simulations

Fully three-dimensional MHD models of the magnetosphere have already been developed for scientific use. Their input parameters are typically solar wind density, velocity, and interplanetary magnetic field. The inner boundary of the magnetosphere is typically set at somewhat above  $3 R_E$ , and physical quantities are mapped down to the ionosphere along field lines. The output is the dynamic response of the magnetosphere-ionosphere system. These models do not generally provide a proper description of the inner magnetosphere, because of (i) the definition of the inner boundary, and (ii) the presence of dominant non-MHD processes in the inner magnetosphere.

This field is very active. Several groups are developing their own model, and have not yet published their results. The situation is then expected to evolve rapidly during the next decade. The best known global MHD models of the magnetosphere are those developed at the University of Maryland (see Mobarry et al., 1996) and at the University of California at Los Angeles (see Raeder et al., 1997 for a recent application). Another model developed at UCLA (Walker et al., 1993) is worth being mentioned here. In Europe, a model is currently being developed for scientific use only at the Finnish Meteorological Institute (Janhunen, 1996).

#### 4.1.3 Kinetic models

The behaviour of the inner magnetosphere cannot be described in the frame of the MHD approximations<sup>4</sup>. Models of the inner magnetosphere should take into account:

- the non-local character of the plasma response for any transport with characteristic time duration longer than the shorter reflection period;
- the effect of the plasma on the magnetic field;
- the interaction between the different waves in the plasma.

Self-consistent kinetic methods properly describe the field-particle interaction and satisfy these requirements at the expense of very heavy calculations. An example of such an approach is the model of substorm growing phases recently developed by Le Contel et al. (1999a, b). Their model is based on the approach proposed by Pellat et al. (1994) and Hurricane et al. (1995) to describe electromagnetic perturbations in a non-adiabatic<sup>5</sup> plasma. It allows one to estimate the adiabatic response of the plasma to a given external electromagnetic perturbation that originates in the solar wind: an east-west current flowing close to the equatorial plane in the far tail.

The computations could become significantly simpler if one neglects the perturbation generated by the motion of particles, which assumes that the magnetic field model used takes into account satisfactorily the existence of the currents associated with the particles whose motion is studied. This approach uses both an Eulerian approach (global model of the magnetosphere) and a Lagrangian one (particle transport). In practice, it provides an efficient tool for describing the trajectories of the particles in the magnetosphere once an accurate model of the magnetospheric magnetic field, for example, Smets (1998) used such an approach to study the particle distribution functions associated with reconnection processes. Numerical simulation of particle transport then allowed him to characterize particle distributions associated inside the magnetosphere to reconnections at sites with different topology and localisation.

<sup>4</sup>The ideal MHD basically assumes that the plasma is a perfectly conductive medium.

<sup>5</sup>Adiabatic responses of a plasma correspond to situations where the motion invariants defined in the frame of the Hamiltonian mechanics are conserved. For a charged particle in the presence of a magnetic field, the first invariant is its magnetic moment  $\mu$ , the second invariant is the integral along the field line of the component along the magnetic field of the particle momentum, and the third one is the magnetic flux encircled by the particle's periodic drift shell orbits.

#### 4.2 Specific models

##### 4.2.1 The Magnetospheric Specification and Forecast Model (MSFM)

The MSFM (Freeman et al., 1994; Lambour et al., 1997) is a large-scale physical model designed to specify fluxes of electrons,  $H^+$ , and  $O^+$  in the energy range responsible for spacecraft charging,  $\sim 100$  eV to  $\sim 100$  keV. It is being developed for operational use by the US Air Force, and it is probably, at present, the closest to being an operational large-scale physical model. Its description is available on the net<sup>6</sup>. It is an update of a series of earlier models. Its predecessor, the MSM (Magnetospheric Specification Model), is routinely used by NOAA/SEC for space weather services. Its major improvement compared to the earlier models, is the complexity of the electric and magnetic field models and its capability to run in real time.

The primary input parameters for MSFM are:

1. the  $K_p$  and  $D_{st}$  geomagnetic indices;
2. the polar cap potential drop and the auroral boundary index that specify the polar ionospheric electric field distribution and the auroral precipitation pattern;
3. the solar wind density and speed, and the interplanetary magnetic field (IMF).

Secondary input parameters include:

1. the sum of  $K_p$ ;
2. precipitating particle flux and polar cap potential profiles from the operational DSMP satellites.

The solar wind density and speed define the magnetopause stand-off distance, and the IMF is used to select the appropriate convection pattern in the polar cap. All together, the primary input parameters determine the electric and magnetic field models used. The sum of  $K_p$  is used as an indicator of the long-term activity level.

The model can operate with reduced sets of input parameters, and in particular, with  $K_p$  alone. It also includes neural network algorithms that predict the input parameters from solar wind measurements. Thus, it has some capability for short-term space weather forecasting.

MSFM follows particle drift through the magnetosphere using slowly time-varying electric and magnetic field models. The electric and magnetic field configurations are updated every 15 min. The particle distribution is isotropic, and the model keeps track of energetic particle loss by charge exchange and electron precipitation into the ionosphere.

MSFM successfully accounts for most major electron flux enhancements observed at geosynchronous orbit. Flux dropouts that often precede the flux enhancements are predicted with less confidence, and they are often missed near the dawn meridian.

<sup>6</sup><http://rigel.rice.edu/~ding/msfm95/msfm.html>

#### 4.2.2 Ring current models

It has already been mentioned that the ring current composition and dynamic particles is driven by the level of the magnetospheric activity. During periods of magnetospheric activity, the ring current dynamics is driven by particle injection: direct ionospheric ion injection, inward transport of plasma sheet and pre-existing ring current particles. On the contrary, the behaviour of the quiescent ring current is driven by diffusion of low energy protons in the presence of convection electric fields (see, e.g. Daglis et al., 1999). Different quantitative models have been developed accordingly for storm- and quiet-time ring currents.

A family of storm-time ring current models neglects direct injection of ions from the ionosphere. The main source of ring current particles is then the plasma sheet, and the characteristics of that source enter the models as boundary conditions. At least four groups are working at present along this path. They are located at the University of Michigan, the Aerospace Corporation, NASA-Goddard Space Flight Centre, and Rice University (see Chen et al., 1994 for a review of these works). Differences between their models are mainly related to a different selection of aspects of the problem that are treated at a state-of-the-art level or neglected. One of the main difficulties for these models in the context of space weather operational use is the specification of the boundary condition (i.e. the distribution function of the particles at all parts of the boundary, where the drift velocity is inward) in absence of adequate real-time data.

More sophisticated simulations have been developed, which involve both global magnetospheric models and Lagrangian models of particle transport. They have been developed to address questions in relation to academic research rather than to space weather activities. Such a coupled investigation is, however, very promising for space weather activities, since it provides an efficient tool for the particle distribution function determination in a given region of the magnetosphere.

With the objective of analysing the ring current dynamics, Fok et al. (1993) have developed a kinetic model that simulates the evolution with time of the distribution functions for the dominant ions ( $H^+$ ,  $He^{++}$ ,  $O^+$ ), as a result of charge exchange and Coulomb collisions. Their results show that during the recovery phase of a magnetic storm, Coulomb collisions lead to generation of low energy ( $< 500$  eV) ions and significant heating of plasmaspheric populations. More recently, the model of Fok et al. (1993) and the particle code of Delcourt et al. (1990a) have been coupled to investigate the ring current response to the field line depolarisation observed in the inner magnetosphere during the substorm growing phase. The code of Delcourt et al. (1990a) computes the path of charged single particles, given time-varying electric and magnetic fields, from their injection in the magnetosphere – from the magnetosheath or the ionosphere – to their input in the plasmasphere.

In a more recent study, Fok et al. (1999a) used the particle code of Delcourt et al. (1990b) to estimate the particle

distribution in the near tail (at a geocentric distance of about  $12 R_E$ ), by means of back propagation of particles to their sources. The transport and acceleration of these ions in the inner magnetosphere, as well as their contribution to the ring current, are then computed using the kinetic model of Fok et al. (1993), with a magnetic field more realistic than a dipolar one. These simulations account for many characteristic features of substorms. Further refinements of these investigations will involve an estimate of the initial particle distribution functions from the results of MHD simulations (see, e.g. Fok et al., 1999b) for given solar wind situations, instead of semi-empirical modelling.

Understanding the loss process in the region where the hot ring current plasma coexists with cold plasmaspheric plasma requires modelling of the interaction between ring current ions and plasma waves resonantly generated by the coexisting hot and cold plasmas. Modelling wave particle interaction on a global scale is very challenging because it needs a self-consistent description of wave and particle behaviour.

#### 4.2.3 Radiation belt models

##### 4.2.3.1 Empirical models

Very early in the space era, it became clear that the average radiation doses received by satellites are key parameters that should be known. The United States and the former Soviet Union, therefore, started building empirical models of radiation belts, using data collected for several years on board a number of satellites.

In the United States, these models were developed for NASA by Aerospace Corporation. The latest versions are the AP8 (Aerospace Protons #8) and AE8 (Aerospace Protons #8) models (Sawyer and Vette, 1976; Vette, 1991). They have been built at the end of the 1970s, using data from about 40 satellites that were in orbit between 1961 and 1977, a period that mostly corresponds to solar cycle #20. They provide proton flux for energy in the range 100 keV to 400 MeV and electron flux for energy in the range 40 keV to 7 MeV, for altitudes up to that of geostationary orbits.

The necessity of updating these models led the United States to launch a satellite dedicated to radiation measurements: the NASA/DoD CRRES (Combined Release and Radiation Effect Satellite) satellite. It has been operational for 14 months – between August 1990 and October 1991 – that roughly corresponds to the maximum of solar cycle # 22. The CRRES data have been used by DoD to develop new empirical models for the proton (CRRESPRO) and electron (CRRESELE) radiation belts (see, e.g. Brautigam et al., 1992; Gussenhoven et al., 1993). These models do not provide a significant updating of earlier models, because (i) the altitude and energy ranges they cover are not as wide as those covered by NASA's AP and AE models, and (ii) they rely on data collected by only one satellite during a disturbed period that is probably not representative of the average behaviour of the radiation belts.

On the contrary, the durability of some meteorological satellites gives the opportunity to improve the models for the low altitude range. Boeing is presently preparing for NASA a model using the NOAA/TIROS satellite data (Huston et al., 1998). This model relies on measurements made since 1978, i.e. over almost two solar cycles. Despite its limitation with respect to altitude ( $< 800$  km) and energy (3 ranges for energy larger than 16 MeV), it is a significant improvement compared to earlier models because it describes the evolution with solar cycle of high energy proton fluxes. The authors are now trying to extend the altitude range addressed by the model.

In addition to the models described above, nucleus of models appear at many places in the world. Most of them rely on data from only one satellite, and their representativeness is limited accordingly. As a result of these limitations, the AP8 and AE8 empirical models will probably still remain the standard to be used for mission planning during the next few years, despite their own limitations.

#### 4.2.3.2 Physical models

Physical models have also been developed very early, in order to elucidate the radiation belt sources and their dynamics. Due to the limitation in the computer capacities, they did not significantly evolve during the 1980s. The increasing efficiency of computers allowed one to restart their development at the beginning of the 1990s. During the last few years, they allowed us to make possible significant breakthroughs in our understanding of the sources and dynamics of the proton and electron belts.

They rely basically on the resolution of the Fokker-Planck diffusion equation in adiabatic conditions. They take into account more and more different physical phenomenon:

- particle-particle interactions;
- wave-particle interactions;
- Coulomb interactions;
- plasma influence.

An example of such code is Salammbô, a set of codes devoted to the understanding of high energy charged particle transport during magnetic activity periods. The first one, called Salammbô-3D, solves the classical Fokker-Planck diffusion equation, either for proton or electron radiation belts, in the 3-D phase space. The equation is classically written in terms of the three adiabatic invariant, corresponding to energy, pitch angle and  $L$  McIlwain parameter, so this version is a real 2-D one in space, where the results are averaged with respect to the longitude (local time) (Beutier et al., 1995; Beutier and Boscher, 1995). The Salammbô-4D version is an extended version of Salammbô-3D, taking into account the longitude (local time), so it is a real 3-D one in space. It was developed to understand injections during substorm periods, and their effects, such as drift echoes or the growth of the

ring current. In this version, it is possible to follow particles in their drift around the Earth (Bourdarie et al., 1997).

Tests made on the prediction of the time dependence of the lifetime of trapped electrons (i.e. the time needed to have the initial flux decreased by a factor  $e$ ) have shown that radiation belt models are very sensitive to the waves (Bausart et al., 2000). The elaboration of a physical radiation belt model, therefore, requires the introduction of a wave model. The two models presently used are the Abel and Thorne semi-empirical model (Abel and Thorne, 1998a, b) and the LPCE/CEA empirical model (Baussart et al., 2000; Lefeuvre et al., 2000) based mainly on statistics performed from 3 years of DE-1 data and on wave normal directions estimated from GEOS-1 and ISEE-1 data.

## 5 The ionosphere-thermosphere system

Above 90 km altitude, the upper thermosphere is characterized by a large density decrease, that implies a decrease in the collision frequency between molecules. The temperature increases rapidly from about 180 K to the thermopause value of about 1000 K. This value, constant above 300 km, is also called the exospheric temperature. It is directly dependent on the solar energy in the UV and EUV bands and on the auroral energy inputs. The main atmospheric constituents, nitrogen and molecular oxygen, are dissociated and ionised by the absorption of EUV solar radiation and particle precipitation to form the ionosphere. The resulting ionised species are mainly molecular ions ( $N_2^+$  and  $O_2^+$ ) at low altitude (below 200 km) and atomic ions ( $O^+$ ) at high altitude (above 200 km).  $N^+$  is created by  $N_2$  dissociative ionisation,  $H^+$  by charge exchange reactions between  $O^+$  and H. Finally,  $NO^+$  arises from chemical recombination with neutrals. It is a major ion of the E-region, although its neutral parent NO is a minor constituent of the atmosphere in that region. The ionosphere can, therefore, be said to be composed of three molecular ions ( $N_2^+$ ,  $NO^+$  and  $O_2^+$ ), three atomic ions ( $H^+$ ,  $N^+$  and  $O^+$ ) and electrons. The altitude of the maximum of the electron density is close to 250 km.

The understanding of the behaviour of this solar-terrestrial sub-system and its effects on technological operations is determined by the ability to model at least the height, geographical and time distributions of the electron concentration and the neutral densities. The modelling effort has started a long time ago, due to the development of our technological society (telecommunications, orbitography, etc.). However, there is no numerical code or model that is able at the present time to describe accurately both the three-dimensional and time-dependent distribution of the ionospheric plasma, and the thermospheric densities during quiet and disturbed conditions. In other words, no model is able to reproduce in a satisfactory way both the climate and the weather of the Earth's ionosphere-thermosphere. In addition, there is no well established experimental database that can be used to verify and test the existing models, in order to generate the improvements needed.

At present, most models are able only to reproduce consistently the climate of the ionosphere and the thermosphere, as defined mostly by diurnal, seasonal and solar cycle variations. Serious theoretical and computational efforts, supported by complex combinations of experimental techniques, are being done to model and predict the weather. This weather is defined as the hour-to-hour, day-to-day, week-to-week variability of the electron and neutral concentrations within the framework of the climatology.

Three different types of models can be identified, each of them having implicit limitations and advantages:

- semi-empirical and empirical models;
- physical models;
- in the case of ionospheric modelling only, analytical “profilers” based on routinely scaled ionospheric data.

The physical models of the ionosphere and the thermosphere use as inputs results from empirical models of the convection electric field and auroral precipitations, in order to describe the coupling with the magnetosphere. Such models are also presented at the end of this section.

## 5.1 Empirical and semi-empirical models

### 5.1.1 Ionospheric models

These models are based on a description of the ionosphere in terms of analytical functions. These functions are estimated either from experimental data or from the results of physical models.

A well-known empirical model widely used for different applications is the International Reference Ionosphere (IRI)<sup>7</sup> (Bilitza, 1990). This is the result of an international project sponsored by the Committee on Space Research (COSPAR) and the International Union of Radio Science (URSI), which aimed at producing a reference model of the ionosphere based on available experimental data sources. For a given location, time and date, IRI describes the electron concentration, electron temperature, ion temperature, and ion composition in the altitude range from about 50 km to about 2000 km, as well as the total electron content (TEC). The solar activity is represented by the sunspot number index. IRI provides monthly averages in the non-auroral ionosphere for magnetically quiet conditions; it can also be used for estimating the profile of electron concentration using the experimental values of F2 peak electron concentration (i.e.  $f_oF2$ ) and height as inputs.

The major IRI data sources are the coefficients ( $f_oF2$  and  $M(3000)$ ) produced by the Radiocommunication Sector of the International Telecommunication Union (ITU-R)<sup>8</sup> on the basis of a large set of ground vertical ionosonde data, the

powerful incoherent scatter radars (Jicamarca, Arecibo, Millstone Hill, Malvern, St. Santin), the ISIS and Alouette topside sounders, and in situ instruments on several satellites and rockets. IRI is updated periodically and has evolved over a number of years.

At present, the major limitation of IRI appears to be its description of the electron distribution in the region above the peak of the F2-region (topside ionosphere). Such distribution gives vertical total electron contents that have values above the expected ones, particularly at middle and high latitudes for high solar activity. The height limit of 2000 km for the electron concentration calculation is a major limit for the use of IRI to provide TEC estimates for satellite heights, such as those of the GPS constellation.

A group of semi-empirical models have been developed by D. N. Anderson and colleagues of the Phillips Laboratory of the USAF. There are based on the combination of databases of coefficients that reproduce theoretically calculated profiles based on physical models:

- the Semi-Empirical Low-Latitude Ionospheric Model (SLIM) (Anderson et al., 1987) is based on a theoretical simulation of the low-latitude ionosphere. Electron concentration profiles are determined for different latitudes and local times by solving the continuity equation for  $O^+$  ions. The profiles are normalised to the F2-peak concentration and are then represented by a Modified Chapman function using six coefficients per individual profiles. Input parameters used in the theoretical calculation include the MSIS model neutral temperatures and densities (see below), the IRI model temperature ratios, and the diurnal ion drift patterns observed by the Jicamarca incoherent scatter radar for the different seasons;
- the Fully Analytical Ionospheric Model (FAIM) (Anderson et al., 1989) uses the formalism of the Chiu model (Chiu, 1975) with coefficients fitted to the SLIM model profiles. The local time variation is expressed by a Fourier series up to order 6, and the variation with dip latitude by a fourth order harmonic oscillator function (Hermite polynomial);
- the same group developed a Parameterised Real-time Ionospheric Specification Model (PRISM) (Daniell et al., 1995) that consists of two segments. One is the Parameterised Ionospheric Model (PIM) and the other incorporates near-real-time data from ground and satellite sensors. PIM is a relatively fast global ionospheric model. It is based on the output of physical ionospheric models developed by the Utah State University and Boston College. It consists of a source code and a large database of runs of ionospheric specification codes. From a given set of geophysical conditions (day of the year, solar activity index  $f_{10.7}$ , geomagnetic activity index  $K_p$  . . .) and positions (latitude, longitude, and altitude), the model can produce critical frequencies and peak heights for the ionospheric E- and

<sup>7</sup>Available on the NSSDC site: [http://nssdc.gsfc.nasa.gov/space/model/models\\_home.html](http://nssdc.gsfc.nasa.gov/space/model/models_home.html)

<sup>8</sup><http://www.itu.int/ITU-R/index.html>

F2-regions, as well as electron concentration profiles from 90 to 1600 km and TEC.

### 5.1.2 Thermospheric models

Thermospheric semi-empirical models are based on the hypothesis of independent static diffusive equilibrium of the different thermospheric constituents above 120 km altitude. The density of each constituent decreases exponentially with altitude, according to its own scale height that depends on its mass and on the temperature. The temperature variation is described by the Bates profile that is a function of the exospheric temperature, and the temperature and its gradient at the lower limit. The temperature and the gas concentration evolve depending on season and solar local time, latitude, solar flux represented by the  $f_{10.7}$  index, and magnetic activity represented by  $A_p$  or  $K_p$  indices. Periodic functions are used to represent the diurnal and seasonal variations. Latitudinal, solar activity and magnetic activity variations are represented by non-periodic functions that may differ from one model to another.

The DTM-78 model (Drag Temperature Model; Barlier et al., 1978) has been developed at the Centre d'Etudes et de Recherches Géodynamiques (CERGA, France) in a co-operative effort with the Institut d'Aéronomie spatiale in Belgium and the Service d'Aéronomie du CNRS in France. It is based on temperature measurements by the OGO-6 satellite and density data from drag observations. The DTM-94 version (Berger et al., 1998) include data from the micro-accelerometer CACTUS (1975–1979) and from the Dynamic Explorer 2 satellite (1981–1983). This version is actually used for various applications: trajectory computation for satellites such as SPOT or TOPEX/POSEIDON and developments of gravity models. A new version is under development by CNES/GRGS and Service d'Aéronomie, in order to better describe the variations in the lowest altitude region (120 to 150 km).

MSIS (Mass Spectrometer and Incoherent Scatter) is the most widely used model in the scientific community. It has been developed by GSFC (Goddard, Maryland) and is based on satellite mass spectrometer data and ground incoherent scatter data. The MSIS-86 version (Hedin, 1987) became the highest part (altitude above 90 km) of the COSPAR International Reference Atmosphere (CIRA). The MSISE-90 version<sup>9</sup> (MSIS Extended, Hedin, 1991) describes the atmosphere from the ground level. Above 72.5 km altitude, it is a revised version of MSIS-86 that includes new data. Above 120 km, MSIS-86 and MSISE-90 are identical.

An empirical model of the thermospheric winds has also been developed by GSFC. The first version, HWM-87 (Horizontal Wind Model), is based on data from the AE-E and DE-2 satellites. Ground-based measurements by incoherent scatter radars and Fabry-Perot interferometers have been added in the version HWM-90 (Hedin et al., 1991) to lower

the altitude limit down to 100 km. The last version HWM-93<sup>10</sup> (Hedin et al., 1996) describes the winds almost from the ground using data from MF and meteor radars.

The MSIS and HWM models are often used as inputs for physical modelling, for example, the ionospheric and radiation belt models. MSISE-90 will be used for orbit prediction of the TIMED spacecraft. Current development efforts have been transferred at NRL (Washington, D. C.) and are focused on the improvement of the solar and geomagnetic input specifications, and on the incorporation of new data sources (Picone et al., 2000).

### 5.2 Physical models

The macroscopic behaviour of the ionospheric molecular ions ( $N_2^+$ ,  $NO^+$  and  $O_2^+$ ) and atomic ions ( $H^+$ ,  $N^+$  and  $O^+$ ) is described through the “fluid approach”. The set of transport equations (continuity, momentum, energy and heat flow) corresponding to this approach is derived from Schunk (1977). It solves the temporal evolution of the concentration, the field-aligned velocity, and the temperature and the field-aligned heat flow of each species.

On the other hand, precipitating electrons or primary photoelectrons move along the magnetic field lines, producing heat, excitation and ionisation. A cascade may occur, producing secondary electrons through electron impact collisions. These microscopic collisions and absorption are described through a kinetic approach. Some outputs of the kinetic equation are the ion and electron productions, and the thermal electron gas heating rate.

Therefore, a physical description of the ionosphere requires the coupling of the two approaches. Global or limited area models that use basic physical principles controlling the ionospheric plasma have been developed in recent years. They make use of empirically specified input data. The main ones are the solar EUV/UV radiation (see Sect. 5.5), the ion convection pattern (see Sect. 5.4) and the auroral precipitation pattern (see Sect. 5.4). The accuracy of these inputs limits the ability of these physical models to reproduce climatic and space weather conditions. One finds in Schunk (1996) an extensive review on the situation up to 1996. In 2002, the main ionospheric models that follow this approach are the Utah State University model of the global ionosphere (Schunk and Sojka, 1996), the University of Alabama Field Line Integrated Plasma (FLIP; Richards and Torr, 1996), the Phillips Laboratory Global Theoretical Ionospheric Model (GTIM; Anderson et al., 1996), and the TRANSCAR model (Lilensten and Blelly, 2002). Zhang et al. (1993) and Zhang and Radicella (1993) developed a time dependent ionospheric model essentially limited to middle latitudes, but with the advantage of being fast in comparison with global model calculations.

All of these models are suitable for the investigation of the influence of different physical processes on the ionosphere.

<sup>9</sup> Available on the NSSDC site: [http://nssdc.gsfc.nasa.gov/space/model/models\\_home.html](http://nssdc.gsfc.nasa.gov/space/model/models_home.html)

<sup>10</sup> Available on the NSSDC site: [http://nssdc.gsfc.nasa.gov/space/model/models\\_home.html](http://nssdc.gsfc.nasa.gov/space/model/models_home.html)

Their main outputs are the electron concentration distribution, electron and ion temperatures, ion velocity, with all parameters measurable by incoherent scatter radar. There are several other outputs, such as electron heating rate, ion and electron production, or ion and neutral excitation rates. From this last parameter, it is possible to compute the electromagnetic emissions due to the deactivation. As an example, Witasse et al. (1999) have added the calculation of the thermospheric red and green oxygen emission lines, in order to use TRANSCAR as a diagnostic tool for the thermosphere.

The physical description of the upper atmosphere requires the need to raise the question of the ionosphere-thermosphere coupling, and possibly the coupling with the lower atmosphere. Coupled approaches necessitate solving self-consistently the 3-D time-dependent equations of momentum, energy and continuity for neutral particles and ions. Parameters relevant to the ion-neutral and neutral-ion coupling are exchanged at every few time steps between the thermospheric and ionospheric codes.

The Coupled Thermosphere-Ionosphere-Plasmasphere or CTIP (Fuller-Rowell et al., 1996; Millward et al., 1996) is one of the most comprehensive upper atmosphere models currently available. It covers the region from 80 km to 450 km altitude in the neutral atmosphere, 120 km to 10 000 km in the ionosphere. This model has been used extensively to study the response of the ionosphere/thermosphere system to geomagnetic storms (Codrescu et al., 1997; Fuller-Rowell et al., 2000). In 1998, it was used in the frame of space weather in an effort to set a network for ionospheric predictions (Szuszczewicz et al., 1998). Different recent code developments have been performed. They include the addition of tidal and planetary wave forcing at the CTIP model's lower boundary (Fuller-Rowell et al., 1999; Muller-Wodarg et al., 2000); self-consistent calculation of the E-region dynamo electric field (Millward et al., 2000); flexible high-latitude auroral precipitation and convection electric field (Schoendorf et al., 1996), including the proton precipitation (Galand et al., 2001); and extension of the lower boundary below the stratopause down to 30 km (Harris, 2000, personal communication).

The NCAR TGCM's are three-dimensional, time-dependent models of the Earth's neutral upper atmosphere (Roble and Ridley, 1994 and references herein). Recent models in the series include a self-consistent aeronomic scheme for the coupled Thermosphere/Ionosphere system, the Thermosphere Ionosphere Electrodynamics General Circulation Model (TIEGCM), and an extension of the lower boundary from 97 to 30 km, including the physical and chemical processes appropriate for the mesosphere and upper stratosphere, the Thermosphere Ionosphere Mesosphere Electrodynamics General Circulation Model (TIME-GCM). The outputs of the TIME-GCM consist of 30 neutral atmospheric parameters on a three-dimensional latitude, longitude, pressure grid. Geographic longitude begins at  $-180^\circ$  west and continues around the globe with a  $5^\circ$  resolution. Geographic latitude resolution is also  $5^\circ$ , from  $-87.5^\circ$  south to  $+87.5^\circ$  north. The vertical dimension is in a log pressure

scale from  $-17.0$  at the bottom (approximately 30 km) to  $5.0$  at the top (varying in altitude up to about 400 km). The applications are numerous, and a very large number of papers have been published in the past years based on the use of this model and comparison with experimental data. Recent applications include a comprehensive study of the atmospheric tides (Hagan and Roble, 2002).

The next step in the physical description of the upper atmosphere is to couple the TIE codes with the upper boundary conditions, i.e. magnetospheric/protonospheric codes. This has been achieved for the first time by the numerical model of the Earth's upper atmosphere (Namgaladze et al., 1991) constructed at the Kaliningrad Observatory of IZMIRAN and at the Polar Geophysical Institute in Murmansk. This model covers the height range from 80 km up to a geocentric distance of 15 Earth radii. It consists of three main blocks: thermospheric, ionospheric-protonospheric and an electric field computation block. The exchange of information between these blocks is carried out at every time step of the numerical integration of the modelling equations (typically 2 min). The electric fields, both of thermospheric dynamo and magnetospheric origin, and protonospheric parameters are calculated consistently. For the thermospheric and molecular ion parameters, the latitudinal integration steps vary from  $10^\circ$  at the geomagnetic equator to  $2^\circ$  at the auroral zones. For the electric field, ionospheric F2-region and protonospheric parameters, they vary, respectively, from  $5^\circ$  to  $2^\circ$ . The longitude step is  $15^\circ$ . The thermosphere is computed between 80 and 520 km at 30 altitude levels. This model has been used in several applications. Recently, the effect of geomagnetic storms on both the thermosphere (Förster et al., 1999) and the ionosphere (Namgaladze et al., 2000) have been investigated.

The Magnetosphere-Thermosphere-Ionosphere-Electrodynamics-General circulation model (MTIE-GCM; Peymirat et al., 1998) couples the TIE-GCM (Richmond et al., 1992) and the Ionosphere-Magnetosphere model of Peymirat and Fontaine (1994). It calculates self-consistently the 3-D structure of the thermosphere and of the ionosphere, the 2-D structure of the magnetospheric plasma convection in the equatorial plane of the magnetosphere, and the couplings between the thermosphere, the ionosphere and the magnetosphere. Once given as inputs, the distribution of the electric potential along the polar cap boundary, and the density and temperature of the magnetospheric plasma source, the magnetospheric part computes the motion of the magnetospheric plasma, from which it deduces the auroral precipitation and the region's two field-aligned currents. These are transmitted to the TIE-GCM part, which computes the 3-D dynamics of the thermosphere and ionosphere. The ionospheric electric field is also computed considering the effect of auroral precipitation, the effect of the region's two field-aligned currents and the effect of the neutral winds. The ionospheric electric field is then transmitted to the magnetospheric part, in order to calculate the motion of the magnetospheric plasma. This model is used to evaluate the feedback coupling among the magnetosphere-ionosphere-

thermosphere system and specifically to check the influence of the neutral winds on this coupling. A coupling of the magnetospheric part to the physical ionospheric code TRANCAR, described above, is under way (Peymirat and Blélly, personal communication, 2001).

### 5.3 Ionospheric profilers

Another way to estimate the electron concentration distribution with height is based on data assimilation techniques. It is based on adjustments of a parameterised profile, expressed in terms of simple mathematical functions, to ionospheric characteristics routinely scaled from the ionograms. The advantage of this type of “profile modelling” is that it can, in principle, use as inputs experimental values of basic ionospheric characteristics for both quiet and disturbed conditions. The so-called “profilers” share with the other types of models the difficulty of describing well the topside ionosphere. This is essentially due to the absence of a well-established experimental database of topside profiles, which is necessary to obtain a clear description of the topside ionosphere behaviour.

The Bradley and Dudeney (1973) “profiler” describes the electron concentration profile up to the peak of the F2-region. It consists of two parabolic layers for the E- and F2-layer and a linear segment in between. This profile generation is still used by the ITU-R HF propagation prediction method. Dudeney (1978) has later proposed a more refined profile that uses routinely scaled characteristics, incorporates combinations of trigonometric-function segments, and provides optional valley and an F1-ledge description. The main advantage of this profile over the previous one is the continuity of the gradient of electron concentration with height across the segment boundaries.

The DGR model, a “profiler” originally introduced by Di Giovanni and Radicella (1990) and improved by Radicella and Zhang (1995), is able to describe the electron concentration profile in the E, F1, F2-regions of the ionosphere by using simple analytical expressions. It is essentially based on the Epstein layer introduced by Rawer (1982) and considers the existence of characteristic points in the profile with coordinates (values of electron concentrations and their height) calculated by means of empirical expressions. The model is constructed as the sum of three Epstein layers that are formally identical. The model calculates the electron concentration profile above the F2 peak, making use of an effective *shape parameter* empirically derived for the topside ionosphere. The total electron content is computed with an analytical expression obtained from the Epstein layers formulation of the model. This model of electron concentration profile has been adopted by the European Commission COST 238 (PRIME) action and it is part of the computer program produced by the action.

A new family of electron concentration “profilers”, which differ in complexity and which have different but related application areas, has been developed (Hochegeger et al., 2000). They are based on the DGR “profiler” concept and also allow for the use of median or instantaneous values, or maps based

on regional or global experimental data. The three models are:

- NeQuick is a quick-run model for ionospheric applications. This model has been adopted in the ionospheric specifications for the European Space Agency EGNOS project;
- COSTprof is a model that can be used for ionospheric and plasmaspheric satellite to ground applications. This model has been adopted by the COST 251 action of the European Commission as the profiler for its electron concentration distribution;
- NeUoG-plas is a model that can be used particularly in assessment studies involving satellite to satellite propagation of radio waves.

The basic input parameters are values  $f_oF2$  and  $M(3000)F2$ . The output of the models is the electron concentration in the ionosphere as a function of height, geographic latitude and longitude, solar activity (given by sunspot number or by 10.7 cm solar radio flux), season (month) and time (Universal Time UT or local time LT). The models also permit one to calculate electron concentration along arbitrarily chosen ray paths and slant, or vertical total electron content up to heights in the plasmasphere, such as those of GPS satellites. The profiles are continuous in all spatial first derivatives (a necessity in applications like ray tracing and location finding).

Above 100 km and up to the F2-layer peak, all three models are identical using a modified DGR formulation. They use the ITU-R coefficients for  $f_oF2$  and  $M(3000)F2$ , and simplified models for  $f_oF1$  and  $f_oE$ , which take into account solar zenith angle, season and solar activity. The topside F-layer for the NeQuick model is again a semi-Epstein layer, but with a height dependent thickness parameter (Radicella and Zhang, 1995). For the COSTprof, the topside ionosphere formulation uses three physical parameters, namely the oxygen scale height at the F2 peak, its height gradient, and the  $O^+ - H^+$  transition height. These three parameters are modelled according to solar activity, season, local geographic latitude, and “modified dip latitude”. NeUoG-plas has an additional geomagnetic field-aligned third part for the “plasmasphere” to model this region in a more realistic way.

### 5.4 Convection electric field and auroral precipitation models

Such empirical models are used as inputs of the physical ionospheric and coupled ionospheric thermospheric 3-D models. Only the most recent and well-known models are presented below.

The Heppner-Maynard-Rich Electric Field Model is a software package<sup>11</sup> that includes several empirical electric convection field models, and the AFGL Precipitation and

<sup>11</sup> Available on the NSSDC site: [http://nssdc.gsfc.nasa.gov/space/model/models\\_home.html](http://nssdc.gsfc.nasa.gov/space/model/models_home.html)



Conductivity Model; the latter for obtaining conductances, currents, and heating. The Heppner-Maynard models are based on OGO 6 and DE 2 electric field measurements and provide the electric potential and field poleward of 60° geomagnetic latitude (Heppner and Maynard, 1987). Seven different models were generated for different Interplanetary Magnetic Field (IMF) conditions. Spherical harmonics to order 11 in magnetic local time and latitude are used in each case. For southward IMF, an explicit variability with geomagnetic activity is included. The Heelis Electric Convection Field Model (Heelis et al., 1982) is included in this package for comparative purposes. Rich and Maynard (1989) illustrate the improvements of their model in relation to the Heelis model and point out the differences to the Millstone Hill Electric Field Model (Holt et al., 1987) in the region of the Harang discontinuity and near the dayside cleft.

Based on the same set of satellite data, the UAF Eulerian Polar Ionosphere Model (Weimer, 1996) is formulated in the corotating frame and has a scaleable horizontal resolution (up to  $100 \times 100$  km). It derives the electric potentials in the high-latitude ionosphere, resulting from any arbitrary combination of the IMF magnitude and orientation, solar wind velocity, and dipole tilt angle. It is currently used to prepare ionospheric convection maps.

The Izmiran Electrodynamic Model (IZMEM, Feldstein et al., 1984) is based on the inversion of geomagnetic ground-based observations. The 1985 version is parameterized by the interplanetary magnetic field (IMF) strength and direction and is available for the three seasons (summer, winter, equinox). Seven parameters of the high-latitude ionospheric electrodynamics can be determined for the specific IMF strength and orientation ( $B_x$ ,  $B_y$ ,  $B_z$ ): geomagnetic perturbation vectors at the Earth's surface, electrostatic field potential at the ionospheric altitude, as well as electric field vectors, field-aligned currents, ionospheric current vectors, equivalent current vectors, and Joule heating rate. The new IZMEM/DMSF<sup>12</sup> model produces patterns obtained after recalibration of the IZMEM model by the observed DMSF electrostatic potentials (Papitashvili and Rich, 2002). It uses real-time IMF data from the ACE satellite.

The AFGL Electron Precipitation Model (Hardy et al., 1987) and the AFGL Ion Precipitation Model (Hardy et al., 1989) provide the integral energy and number flux of precipitating auroral electrons and ions as a function of corrected geomagnetic latitude, magnetic local time, and magnetic activity ( $K_p$ ). These are based on millions of spectra from the DMSF-F2, F4, F6, and F7 satellites and the P78-1 satellite. At each level of activity, the high-latitude region was separated into 30 zones in corrected geomagnetic latitude (from 50° to 90°) and 48 one-half-hour zones in magnetic local time. The electron model also provides Pedersen and Hall conductivities, using empirical relationships between conductivities, electron energy flux, and average energy.

For space weather applications, it should be noted that real-time global convection maps, or the equivalent electro-

static potential maps, constrained by SuperDARN radar measurements, have recently been made available<sup>13</sup> (Shepherd and Rhuohoniemi, 2000). Convection pattern and the auroral precipitation pattern can also be deduced from all available data (geomagnetic and radar ground-based observations and in situ measurement of electric field and auroral precipitation) using the AMIE procedure (Richmond, 1992) which today is not a real-time procedure. The use of such time dependant auroral inputs has dramatically increased the physical model's reliability for describing the ionosphere and thermosphere during disturbed conditions.

### 5.5 EUV/UV models for aeronomy

The ultraviolet (UV)/extreme ultraviolet (EUV) solar flux is energetic enough to ionize the upper atmosphere. It constitutes the major source for the diurnal ionosphere. Most of the current models rely on a few experiments taken on board the Dynamics Explorer missions (Hinteregger et al., 1973). A first representation of solar EUV fluxes for aeronomical applications was given by Hinteregger (1981), and Hinteregger and Katsura (1981). A first reference flux SC#21REF was assembled from measurements performed in July 1976, at a period where  $f_{10.7} = 70$ , and given for 1659 wavelengths. An extrapolation model (SERF 1) allows one to estimate the flux during other periods of solar activity.

Torr and Torr (1979, 1985) proposed two reference fluxes for aeronomy called F79050N – corresponding to  $f_{10.7} = 243$  – and SC#REFW – corresponding to  $f_{10.7} = 68$ . The UV spectrum was divided in 37 bins. Some bins correspond to intense spectral lines, but the bright Lyman  $\alpha$  line at 121.565 nm does not show up because it is not energetic enough to ionise the terrestrial ionosphere. This work proved to be extremely useful. One of its qualities was that the authors proposed the corresponding absorption and ionisation cross sections for the major thermospheric species. Its main limitation is that the measurements do not allow one to reach a good estimate of the flux variability for different solar activity conditions.

Since then, several authors developed their codes in order to take better advantage of the AE database. Amongst them, two are worth noting here:

- Tobiska (Tobiska, 1991; Tobiska and Eparvier, 1998) developed a model called EUV, which takes data from other sources into account: SME, OSO; AEROS; rockets and ground-based facilities. This model takes into account the solar emission zone of each line, through a parameter. It proposes a formula to retrieve a solar flux from the gift of the decimetric index and its average. A new version, SOLAR2000, has been recently developed. It uses a new input parameter named E10.7, computed from a previous version of the code (Tobiska et al., 2000);
- The second improved model is EUVAC (Richards et al., 1994). Its main difference with previous models is

<sup>12</sup><http://www.sprl.umich.edu/MIST/spw.html>

<sup>13</sup><http://superdarn.jhuapl.edu/index.html>

**Table 2.** Synthesis of model inputs and outputs of pre-operational and operational models. Only the ones most often used in a given category are indicated. A column indicates if the model is a physical one (*P*) or a semi-empirical or empirical one (*E*) and 2 columns state whether development of the model itself or availability of new observations is the most crucial for its improvement

	Models	Based on	Inputs	Outputs	<i>P</i> or <i>E</i>	Model develop- ment needed	Observations needed for improving the model	Remark
Sun	Torr	DE satellite data	F <sub>10.7</sub>	EUV/UV spectrum	<i>E</i>		X	New models under construction
Magne- sphere	Shue et al.	Satellite data	IMF Bz Solar wind density and velocity	magnetopause size and shape	<i>E</i>		X	
	Tsyganenko	IGRF Satellite magnetic field data	Geomagnetic indices IMF Solar wind density and velocity	Magnetic Field	<i>E</i>			
	MSM	IGRF	Geomagnetic indices IMF Solar wind density and velocity	Magnetospheric particle fluxes: 0.1–100 Kev Magnetospheric E and B Precipitating electrons Ionospheric convection	<i>P</i>			Next generation: MSFM
	AP8-AE8	Satellite data from 1960–1975	Solar activity: min or max	Particle fluxes Protons: 0.1–400 MeV Electrons: 0.1–7 MeV	<i>E</i>		X	Magnetic storm effects non included
	LPCE/CEA Wave model	Satellite data: DE-1	Geomagnetic indices	Magnetic field power density spectrum Wave distribution Function	<i>E</i>		X	Presently only 2 Kp classes
Iono- sphere	Heppner- Maynard-Rich	Satellite data	Geomagnetic indices IMF	Convection electric field	<i>E</i>			
	SuperDARN	Radar data assimilation	IMF SuperDARN radar data	Convection electric field	<i>E</i>		X	Real time
	AFGL precipitation models	Satellite data	Geomagnetic indices	Ion and electron auroral precipitation Conductivities	<i>E</i>			
	IRI	Ionosondes, IS data Topside sounders	F <sub>10.7</sub> or Sunspot number [NmF2 and HmF2, or TEC]	Electron and ion densities and temperatures	<i>E</i>	X	Mainly topside densities	Valid in non auroral zone and for quiet conditions
	SLIM-FAIM PIM	outputs of physical models	F <sub>10.7</sub> Geomagnetic indices	Electron density profiles	<i>E</i>	X		
	Profilers (DGR derived)	Data assimilation (Ionosondes)	F <sub>10.7</sub> or Sunspot number Ionosondes data	Electron densities profiles	<i>E</i>	X	Mainly topside density	Development co-ordination by COST271 and ESA
Thermo- sphere	MSIS DTM	Satellite data from 1975–1983 + IS data for MSIS	F <sub>10.7</sub> Geomagnetic indices	Neutral densities and temperature	<i>E</i>		X	Mostly valid at mid and low latitudes
	HWM	Satellite data from 1975–1983 IS and inter- ferometer data	F <sub>10.7</sub> Geomagnetic indices	Neutral wind	<i>E</i>		X	Mostly valid at mid and low latitudes

the reference flux chosen, and the interpolation formula. The coronal flux is also constrained to be at most 80% of the total.

Those models are very important for aeronomic computation. They allow one to develop fairly good physics. How-

ever, they cannot take into account the variability at different wavelengths. Indeed, specific full measurements (for example, on board the space shuttle) clearly show that there is no linear variation with the decimetric index. On the contrary, some lines of the solar flux may increase with solar activity

**Table 3.** Summary of non-operational physical model classes. The last column states their operational potentialities as foreseen today, if resources are available for operational development

	Models	Method	Inputs	Outputs	Operational potentialities
Sun	Solar atmosphere	MHD	Photospheric magnetic field	Coronal magnetic field; Eruptive phenomena	Long term
Interplanetary medium	Solar Wind	MHD	Plasma and magnetic field at the Sun	IMF, Plasma velocity, density and temperature at 1 AU	Long term
	Global Structure propagation of CMEs	MHD	Plasma and magnetic field at the Sun	Properties of CMEs at 1 AU and beyond	Long term
Magnetosphere	Global models	MHD and/or Kinetic	IMF, solar wind velocity and density at 1 AU Present status of the magnetosphere	Magnetic field and particle distribution in the magnetosphere	Depend on the model
	Radiation belt	Adiabatic invariants Fokker Planck diffusion equation	Geomagnetic indices Thermospheric neutral densities Cosmic neutron flux, Waves Plasmaphere position	Particle fluxes Protons: 0.1–300 MeV Electrons: 0.1–10 MeV	Short term for Salammbô
Ionosphere Thermosphere	Ionosphere 1D and 3D models	Fluid description (+Kinetic Transport)	$F_{10.7}$ Thermospheric neutral densities, temperature and winds Geomagnetic indices or Convection pattern and Auroral Precipitation	Electron and ion densities and temperatures Ion velocity	Short term
	3D Coupled Ionosphere Thermosphere	Fluid description Eulerian approach	$F_{10.7}$ Geomagnetic indices or Convection pattern and Auroral Precipitation	Neutral densities, temperature and wind	Mid term

or decrease, sometimes drastically.

Finally, a radically different approach has been considered by Warren et al. (1996, 1998). They combined a spectral emission line database, solar emission distributions, and estimates from ground-based solar images of the fraction of the Sun covered by the various types of activity, to synthesise the irradiance. The aim was to provide a way to estimate the irradiance by means of the model, from the EUV line emission formed in the upper chromosphere, and in the lower transition region from the Ca II K-line. This approach made it possible for them to estimate the emission measure from a spectrum of a portion of the quiet solar disk measured with the Harvard instrument on Skylab, and compilations of atomic data (Warren et al., 1998). The irradiance spectrum from 50–1200 Å was then computed for the quiet Sun, with the contributions of optically thick emission lines and continua included empirically. A comparison with the empirical models described above indicates relatively good agreement among fluxes of emission lines formed in the solar chromosphere and transition region. A factor of typically 2 is found with the fluxes for coronal emission lines.

## 6 Model synthesis

Tables 2 and 3 summarise the main models described in Sect. 1, their inputs and their outputs. Table 2 presents the

pre-operational and operational models. It is not restricted to the models that are effectively used in operational centres, but includes those that could be used operationally today, as are, for example, most of the empirical models. This table does not, however, include all the models described in the previous sections, but is limited to either recent models or the ones most often used for the same outputs. A letter indicates if the model is a physical one (*P*) or a semi-empirical or empirical one (*E*). We have also indicated whether development of the model itself or availability of new observations is the most crucial for its improvement. The type of data for which the model is based upon is also indicated (column 3), and finally, some remarks are added (column 4). Most of these models use the geomagnetic indices as inputs. The SuperDARN convection model and the ionospheric profilers are based on data assimilation techniques, in order to provide a 3-D description of a given parameter using localised observations as inputs. It should be noted here that new methods are now developed for extracting thermospheric parameters (neutral composition, temperature, total density, winds) from routine ionospheric observations (ground-based vertical sounding, TEC data).

Table 3 presents a summary of non-operational physical models. Contrarily to Table 2, this table deals with classes of models and not with specific models. All models require either improvements or new observations – including for val-

**Table A1:** Members of the consortium lead by ALCATEL Space and LPCE

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(1): contact persons

idation purposes – or both. Most of them are today scientific codes developed for research purposes. The last column states their operational potentialities as foreseen today, if resources are available for operational development. Ring current models, as defined in Sect. 1, are not included in this table, because the data needed at the boundary make it almost impossible to use for operational space weather applications.

## 7 Conclusion

The large variability of the solar terrestrial system limits the efficiency of empirical models, while progress in space weather requires the development of models that rely on basic physical principles controlling the behaviour of the Sun–Earth interaction. Achieving the development of such quantitative physical models is one of the main challenges for space weather issues, and more generally, for solar terrestrial physics. However, one can foresee a long time before this is achieved. The short overview of the state of the art presented in the previous sections shows that this delay is highly dependant on the considered sub-system.

Space weather aims at routinely producing forecasts of relevant parameters. Before the development of operational versions of physical models, this will be achieved by means of hybrid models involving simultaneously physical codes, artificial intelligence-based codes, and empirical models. Such a hybrid approach is presently, and likely to remain for a long time the most promising way of development of operational

models. More developments and improvements of empirical models remain, therefore, mandatory.

These developments imply routine availability of relevant direct measurements or scientifically agreed proxies. In many cases, homogeneous and long enough data series are necessary to assess the models in the various situations that can be encountered in solar terrestrial physics. This is already the case for some ground-based measurements, for example, ionosonde measurements and in space for the GPS satellite data over the last few years. However, it is not yet the case for most in situ measurements on board satellites, inside or outside the magnetosphere. In addition, the indices monitoring solar and geomagnetic activity have not, at present, the required characteristics to comply with the requirements of space weather: for example, time and/or space resolution are not good enough (magnetic indices), measured quantity is not well suited (solar decimetric index). The development of new indices with a better time and/or space resolution, or based on more relevant quantities has already been initiated in the scientific community.

*Acknowledgements.* We thank the Consortium team for fruitful scientific discussions. We also thank P. Cargill for specific input. This work has been initiated in the frame of the preparation of a workshop of the French “Programme National Soleil-Terre”.

The Editor in Chief thanks R. Schwenn and another referee for their help in evaluating this paper.

## Annex 1:

### ESA study for a space weather program – ALCATEL driven consortium

Areas addressed within this ESA study include:

- the user's requirements assessment and market analysis; the parameters needed for space weather and their associated measurement requirements;
- the combined space segment/ground segment scenarios and hardware assessment for meeting the specific needs of an operational space weather system;
- the prediction service definition and its prototyping, including models evaluation;
- the programme development assessment with its European organisation.

The main outcomes expected are comprehensive answers to the following fundamental questions:

- For which user shall a space weather service be developed and in which time frame?
- Which components shall be deployed and with which complexity?
- When shall/can a space weather service be operational?
- How shall a space weather service be implemented, i.e. with which organisation and coordination within Europe and which level of autonomy?

For any further information concerning the tasks performed by the consortium, the reader may refer to the following web site: <http://www.european-space-weather-at.com>.

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